

Unsteady MHD Mixed Convection Flow, Heat and Mass Transfer over an Exponentially Stretching Sheet with Suction, Thermal Radiation and Hall Effect

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Abstract: In this paper the characteristics of flow, radiative heat and mass transfer of a viscous fluid flow over a permeable sheet stretching exponentially with Hall currents in the presence of heat generation and first order chemical reaction are investigated. The non – linear partial differential equations governing the flow are transformed into a set of self similar equations which are then solved numerically. The computational results are graphically presented and discussed. The analysis shows that the Lorentz force resists the primary velocity and the Hall parameter has an opposite effect on it while secondary velocity experiences a reversal effect. The radiation parameter and magnetic field produce thicker thermal boundary layers.

Keywords: Exponentially Stretching Sheet, Hall Currents, MHD, Unsteady.

I. Introduction

Magyari and Keller [1] are the first researchers to investigate the free convection flow, heat and mass transfer of an incompressible fluid from an exponentially stretching vertical surface. They obtained similarity solutions pertaining to an exponential stretching and exponential temperature of the continuous surface. Following the study of Magyari and Keller [1], Elbashbeshy [2] discussed the heat transfer over an exponentially stretching surface with suction. Partha et al. [3] investigated the flow past an exponentially stretching surface of a viscous fluid taking into account of viscous heating. This study was extended by Sajid and Hayat [4] by considering the thermal radiation and an analytical solution was obtained using homotopy analysis method. Later Bidin and Nazar [5] made a numerical study of this problem. Sanjayanand and Khan [6] explored the heat and mass transfer characteristics of a viscoelastic fluid flow due to a stretching surface with viscous dissipation and elastic deformation. Ishak [7] analysed the MHD boundary layer flow and heat transfer of a viscous fluid over an exponentially stretching sheet taking the effect of thermal radiation. Nadeem et al. [8] examined the effect of thermal radiation on the boundary layer flow of a Jeffery fluid induced by a surface stretching exponentially. Battacharyya [9] studied the boundary layer flow and heat transfer over an exponentially sheet. Swati Mukhopadyay et al. [10] studied the features of flow and heat transfer of a viscous incompressible fluid past a permeable exponential stretching sheet with thermal radiation effect. Elbashbeshy et al. [11] made a numerical investigation of flow and heat transfer characteristics of an incompressible fluid over an exponentially stretching surface with thermal radiation. Mukhopadyay et al. [12] investigated the influence of a chemically reactive solute and velocity slip on the boundary layer flow and mass transfer towards an exponentially permeable stretching plate. Nadeem and Hussain [13] investigated the heat transfer of a pseudo plastic fluid past an exponentially porous stretching surface modeling as a Williamson fluid. The steady MHD mixed convective flow past an exponentially stretching sheet was examined by Aziz and Nabil [14] taking thermal radiation and Hall currents.

In this paper we made an attempt to analyse influence of thermal radiation, temperature dependent heat source on the unsteady boundary layer flow heat and mass transfer of a viscous fluid over an exponentially stretching sheet with Hall currents and first order chemical reaction.

II. Mathematical formulation

Consider an unsteady three-dimensional mixed convection boundary layer flow of an incompressible viscous fluid along a stretching surface. The x-axis is taken along the stretching surface in the direction of motion and the y-axis is perpendicular to it. The stretching surface has the velocity $U_w(x, t) = U_0(1 - \alpha t)^{-1}e^{x/L}$, the temperature distribution $T_w(x, t) = T_\infty + T_0(1 - \alpha t)^{-2}e^{x/2L}$ and the concentration distribution $C_w(x, t) = C_\infty + C_0(1 - \alpha t)^{-2}e^{x/2L}$ where U_0 is the reference velocity, α is a positive constant with dimension reciprocal time, L is the reference length, t is the time, T_∞ is the fluid temperature far away from the stretching surface, T_0 is the fluid temperature adjacent to the stretching surface, C_∞ is the fluid concentration far away from the stretching surface and C_0 is the fluid concentration adjacent to the stretching surface. A uniform magnetic field of strength B_0 is applied normally to the stretching surface which produces magnetic effect in the x-axis. The effect of the induced magnetic field is neglected by taking a small magnetic Reynolds number. The

continuity, momentum, energy and concentration equations governing such type of flow invoking the Boussinesq's approximation can be written as

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + g\beta(T - T_\infty) + g\beta^*(C - C_\infty) - \frac{\sigma B_0^2}{\rho(1+m^2)}(u + mw) \tag{2}$$

$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} = \nu \frac{\partial^2 w}{\partial y^2} + \frac{\sigma B_0^2}{\rho(1+m^2)}(mu - w) \tag{3}$$

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{K}{\rho c_p} \frac{\partial^2 T}{\partial y^2} - \frac{1}{\rho c_p} \frac{\partial q_r}{\partial y} + \frac{Q_0}{\rho c_p} (T - T_\infty) \tag{4}$$

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D \frac{\partial^2 C}{\partial y^2} - k(C - C_\infty) \tag{5}$$

The radiative heat flux by using Rosseland approximation can be written as

$$q_r = -\frac{4\sigma_s}{3k^*} \frac{\partial T^4}{\partial y} \tag{6}$$

where σ_s is the Stefan-Boltzmann constant and k^* is the absorption coefficient.

T^4 may be linearly expanded in a Taylor's series about T_∞ to get

$$T^4 = T_\infty^4 + 4T_\infty^3(T - T_\infty) + 6T_\infty^2(T - T_\infty)^2 + \dots,$$

and neglecting higher order terms beyond the first degree in $(T - T_\infty)$,

$$\text{we obtain } T^4 \cong 4T_\infty^3 T - 3T_\infty^4$$

$$q_r = -\frac{16\sigma_s T_\infty^3}{3k^*} \frac{\partial^2 T}{\partial y^2} \tag{7}$$

Substituting equation (7) into equation (4) to get

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{1}{\rho c_p} \left[\left(K + \frac{16\sigma_s T_\infty^3}{3k^*} \right) \frac{\partial^2 T}{\partial y^2} + Q_0(T - T_\infty) \right] \tag{8}$$

The boundary conditions are

$$\begin{aligned} u &= U_w(x, t), v = -V_w(x, t), w = 0, T = T_w(x, t), C = C_w(x, t) \text{ at } y = 0, \\ u &\rightarrow 0, w \rightarrow 0, T \rightarrow T_\infty, C \rightarrow C_\infty \text{ as } y \rightarrow \infty \end{aligned} \tag{9}$$

where u, v and w are the fluid velocity components along x, y and z axes, respectively, ν is the kinematic viscosity, g is the gravity field, β is the volumetric coefficient of thermal expansion, β^* is the coefficient of expansion with concentration, ρ is the density of the fluid, T is the fluid temperature, C is the fluid concentration, σ is the electrical conductivity, c_p is the specific heat at constant pressure, K is the thermal conductivity of the medium, q_r is the radiation heat flux, Q_0 is the uniform volumetric heat generation and absorption, D is the mass diffusivity, k is the chemical reaction, $V_w(x, t) = f_w(U_0\nu/2L(1-\alpha t))^{1/2}e^{x/2L}$ is the velocity of suction ($V_w > 0$), $f_w \geq 0$ is the suction parameter.

We introduce the stream function $\psi(x, y)$ such that $u = \frac{\partial \psi}{\partial y}$ and $v = -\frac{\partial \psi}{\partial x}$.

III. Method of Solution

The governing partial differential equations (2), (3), (5) and (8) can be reduced to a set of ordinary differential equations on introducing the following similarity variables:

$$\eta = \sqrt{\frac{U_0}{2\nu L(1-\alpha t)}} e^{x/2L} y, \tag{10}$$

$$w = \frac{U_0}{(1-\alpha t)} e^{x/L} h(\eta), \tag{11}$$

$$\psi(x, y) = \sqrt{\frac{2U_0\nu L}{(1-\alpha t)}} e^{x/2L} f(\eta), \tag{12}$$

$$T = T_\infty + \frac{T_0}{(1-\alpha t)^2} e^{x/2L} \theta(\eta), \tag{13}$$

$$C = C_\infty + \frac{C_0}{(1-\alpha t)^2} e^{x/2L} \phi(\eta), \tag{14}$$

Using (10) to (14) in equations (2), (3), (5) and (8) we obtain the following set of ordinary differential equations:

$$f''' + ff'' - 2f'^2 - Le^{-X} \left(A(2f' + \eta f'') + \frac{2M}{1+m^2} (f' + mh) - 2e^{-X/2} (G_r\theta + G_c\phi) \right) = 0 \tag{15}$$

$$h'' + fh' - 2f'h - Le^{-X} \left(A(2h + \eta h') - \frac{2M}{1+m^2} (mf' - h) \right) = 0 \tag{16}$$

$$(1 + Nr)\theta'' + Pr(f\theta' - f'\theta - Le^{-X}A(4\theta + \eta\theta')) + \delta\theta = 0 \tag{17}$$

$$\phi'' + Sc(f\phi' - f'\phi - Le^{-X}A(4\phi + \eta\phi')) - \gamma\phi = 0 \tag{18}$$

Where the primes denote the differentiation with respect to η , $X = x/L$ is the dimensionless coordinate, $A = \alpha/U_0$ is the unsteadiness parameter, $M = \sigma B_0^2(1-\alpha t)/U_0\rho$ is the Magnetic parameter, $G_r = g\beta T_0/U_0^2$ is the thermal Grashof number, $G_c = g\beta^*C_0/U_0^2$ is the solutal Grashof number, $Pr = \rho c_p \nu/K$ is

the Prandtl number, $Nr = 16\sigma_s T_\infty^3 / 3Kk^*$ is the thermal radiation parameter, $\delta = 2Q_0L / U_w \rho c_p$ is the heat generation ($\delta > 0$) and absorption ($\delta < 0$) parameter, $Sc = \nu / D$ is the Schmidt number and $\gamma = 2kL / U_w$ is the chemical reaction parameter.

The corresponding boundary conditions are

$$\eta = 0 : f = f_w, f' = 1, h = 0, \theta = 1, \phi = 1, \tag{19}$$

$$\eta \rightarrow \infty : f' \rightarrow 0, h \rightarrow 0, \theta \rightarrow 0, \phi \rightarrow 0, \tag{20}$$

The physical quantities of engineering interest in this problem are the skin friction coefficient C_f and the local Nusselt number Nu_x and local Sherwood number Sh_x which are defined as

$$C_{fx} = \frac{2\mu(\partial u / \partial y)_{y=0}}{\sqrt{Re_x}}, \quad C_{fz} = \frac{2\mu(\partial w / \partial y)_{y=0}}{\sqrt{Re_x}},$$

$$Nu_x = -\frac{x(\partial T / \partial y)_{y=0}}{T_w - T_\infty}, \quad Sh_x = -\frac{x(\partial C / \partial y)_{y=0}}{C_w - C_\infty}, \tag{21}$$

$$\frac{1}{2} C_{fx} \sqrt{Re_x} = f''(0), \quad \frac{1}{2} C_{fz} \sqrt{Re_x} = h'(0),$$

$$Nu_x / \sqrt{Re_x} = -\theta'(0), \quad Sh_x / \sqrt{Re_x} = -\phi'(0), \tag{22}$$

where $\mu = k / \rho c_p$ is the dynamic viscosity of the fluid and $Re_x = x U_w / \nu$ is Reynolds number.

IV. Results and Discussion

The present analysis aims at analysing how the flow characteristics are influenced by Hall currents, chemical reaction, temperature dependent heat source and thermal radiation. Runge – Kutta – Fehlberg method is employed to obtain the numerical solution of equations (15)-(20). To validate numerical solution, the temperature gradient on the surface of the present study is compared with that evaluated by Elbashbeshy et al. [11] neglecting Hall currents and mass diffusion in the unsteady case and Ishak [7] in the steady case and with Elbashbeshy [15], Bidin and Nazar [5] ignoring Lorentz force, thermal and solutal buoyancy, heat generation/absorption and suction. $X = A = Gr = Gc = \delta = M = f_w = \phi = 0$ for different values of Pr and Nr. The compared values of $-\theta'(0)$ are presented in Table 1 which are found to be in good agreement with the said published results.

Table 1 Comparison values of $-\theta'(0)$ for $X=A=\lambda=\delta=f_w=\phi=0$

Nr	M	Pr	Elbashbeshy [15]	Bidin and Nazar [5]	Ishak [7]	Elbashbeshy et al. [11]	Present results	
0.0	0.0	0.72	0.76778	0.9548	0.9548	0.76728	0.76728	
		1.0	0.95478			0.95478	0.95478	
		2.0				1.4715	1.47146	
		3.0	1.86907			1.8691	1.86907	1.86907
		5.0				2.5001	2.50013	2.50013
		10.0	3.66037			3.6604	3.66037	3.66036
0.0	0.5	1.0			0.8611	0.86109	0.86109	
2/3		1.0		0.6765		0.67650	0.67650	
		3.0		1.3807		1.38075	1.38075	

In order to analyze the effects of Hall current, thermal buoyancy force, solutal buoyancy force, chemical reaction, heat source/absorption, thermal radiation and time on the flow-field, the computational values of the primary and secondary velocities of the fluid in the boundary layer region, temperature distribution and mass concentration are presented graphically for various values of magnetic parameter M, Hall current parameter m, suction parameter f_w , thermal Grashof number Gr, solutal Grashof number Gc, Prandtl number Pr, unsteady parameter A, thermal radiation parameter Nr, chemical reaction parameter γ and Schmidt number Sc.

Fig. 1 and Fig. 2 reveal that the primary velocity decreases throughout the boundary layer region with increasing values of magnetic parameter. The deceleration in the velocity is owing to the retarding action of the Lorentz force. Secondary velocity increases rapidly near the plate attaining maximum value and then decreases in the rest of the region eventually approaching the free stream value. The secondary velocity significantly increases with the magnetic field unlike the primary velocity. Fig. 3 and Fig. 4 show the influence of Hall current on the primary velocity and secondary velocity. It is observed that on increasing Hall parameter (m) the primary velocity increases nominally in the vicinity of the plate while the secondary velocity increases significantly throughout boundary layer region. This shows that Hall current tends to accelerate the fluid throughout the boundary layer region which is in conformity of the fact that the secondary velocity arises due to Hall currents. Figs. 5 – 8 illustrate the variation of velocities, temperature and concentration with unsteady parameter. It is revealed that the thickness of the hydrodynamic boundary layer, thermal boundary layer and the solutal boundary layer decreases with increase in the unsteady parameter. Figs. 9 – 12 depict the effect of thermal and solutal buoyancy forces on the primary and secondary velocities. It is observed that both the

velocities increase on increasing Gr and Gc. As Gr (Gc) amounts to the relative strength of thermal (solotal) buoyancy force to viscous force, both thermal and concentration buoyancy forces tend to accelerate both the velocities throughout the boundary layer region. The influence of suction/injection f_w on the velocity distributions are plotted in Fig. 13 and Fig.14. It is observed that primary velocity decreases with increasing suction parameter and the injection (blowing) accelerates the flow. The wall suction($f_w > 0$) results in thinner boundary layers with a fall in the velocity. For blowing ($f_w < 0$) an opposite trend is noticed. The behavior of the secondary velocity with f_w is similar to that of the primary velocity. However, in this case the influence of f_w is stronger than that on primary velocity. $f_w = 0$ correspond to the impermeable stretching surface.

Fig. 15 indicates that increasing values of Prandtl number decrease the temperature. The reduction in the thickness of the thermal boundary layer is justified in view of the fact that increasing values of Pr correspond to reduction in the thermal conductivity resulting in decrease in temperature. The influence of thermal radiation on temperature is depicted in Fig. 16. It is observed that the temperature is enhanced for increasing values of Nr and thus resulting in the increasing of the thickness of both thermal and hydrodynamic boundary layers. The thermal radiation facilitates additional means to diffuse energy as an enhancement in the radiation parameter which corresponds to a reduction in the Rosseland mean absorption coefficient K^* for fixed values of T_∞ and k. Fig. 17 reveals that the temperature is enhanced in the presence of heat source. The heat source releases energy in the thermal boundary layer resulting in the rise of temperature. On increasing $\delta > 0$ the temperature further rises. In the case of heat absorption $\delta < 0$ (heat sink) the temperature falls with decreasing values of $\delta < 0$ owing to the absorption of energy in the thermal boundary layer. From Fig. 18 the species concentration is observed to reduce with increasing values of the Schmidt number throughout the region which is associated with thinner solotal boundary layers. Physically, increasing values of Sc imply decrease of molecular diffusion D. Thus the mass diffusion leads to an enhancement in the species concentration. The influence of chemical reaction rate parameter γ on the species concentration for generative chemical reaction is depicted in Fig. 19. It is found that species concentration with its highest value at the plate decreases slowly till it reaches the minimum value i.e., zero at the far downstream. Further, increasing of the chemical reaction decreases concentration of species in the boundary layer due to the fact that destructive chemical reduces the thickness of the solotal boundary layer and increases the mass transfer. The skin friction coefficient, local Nusselt number and local Sherwood number for different values of the governing parameters are presented in Table 2. The skin friction coefficient in the x-direction reduces for increasing values of the unsteady parameter while it increases in the z-direction. The rate of heat transfer is observed to increase with increasing unsteady parameter while the mass transfer is more for small times. The increase in the magnetic parameter lowers the skin friction coefficient corresponding to the primary velocity and while skin friction coefficient corresponding to the secondary velocity, temperature and the Sherwood number show an increasing tendency. The influence of the Hall parameter is to increase the skin friction coefficient in both x and z – directions, Nusselt number and Sherwood number. The thermal buoyancy and solotal buoyancy parameters both enhance the skin friction coefficient in the x and z – directions, the Nusselt and Sherwood numbers. The Prandtl number depreciates the skin friction coefficient in the x and z-directions and the Sherwood number whereas the Nusselt number shows an increasing tendency. The Nusselt number is seen to decrease with increasing radiation parameter while the skin friction coefficient in both the directions and Sherwood number are found to increase. The influence of Schmidt number on the skin friction coefficient in both the directions is to enhance. The Nusselt number reduces while the Sherwood number enhances with increasing Schmidt number. The increasing value of the heat generation parameter enhances the skin friction coefficient in the x and z – directions and the Sherwood number while the Nusselt number reduces. The chemical reaction parameter decreases the skin friction coefficient and Nusselt number whereas the Sherwood number enhances. The suction parameter reduces the skin friction coefficient in the x-direction while it increases the skin friction coefficient in the z – direction. The Nusselt number and Sherwood number are seen to increase with f_w .

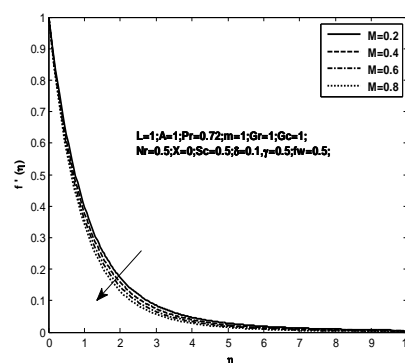


Fig. 1: Primary velocity profiles for different Values of M

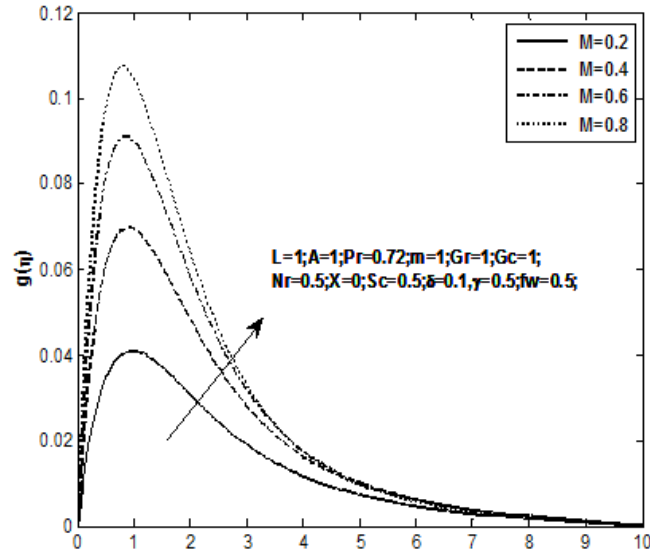


Fig. 2: Secondary velocity profiles for different Values of M

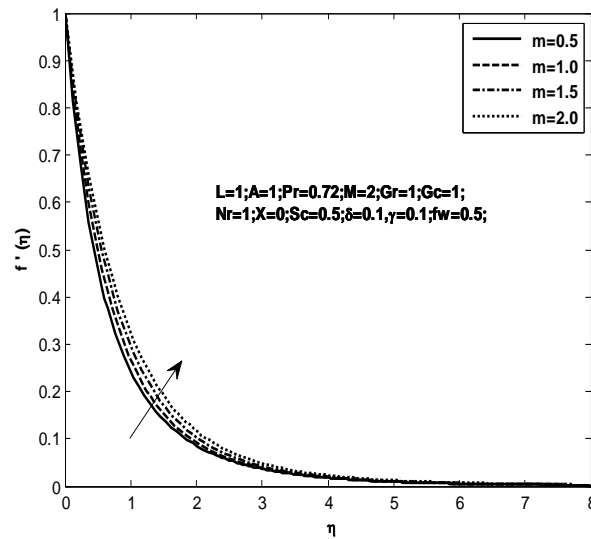


Fig. 3: Primary velocity profiles for different values of m

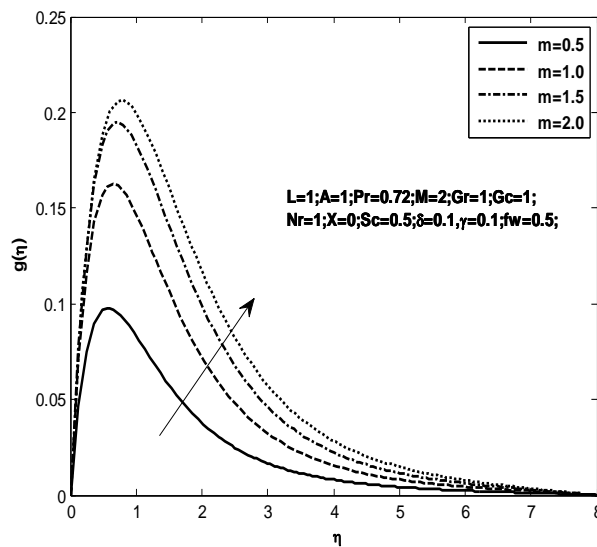


Fig. 4: Secondary velocity profiles for different values of m

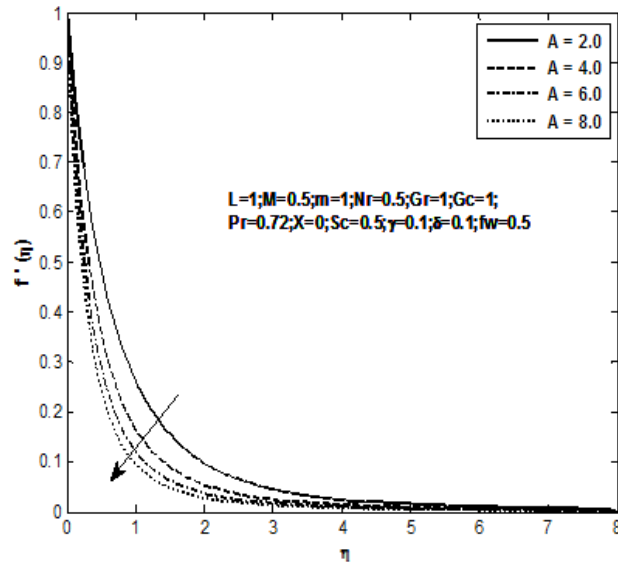


Fig. 5: Primary velocity profiles for different values of A

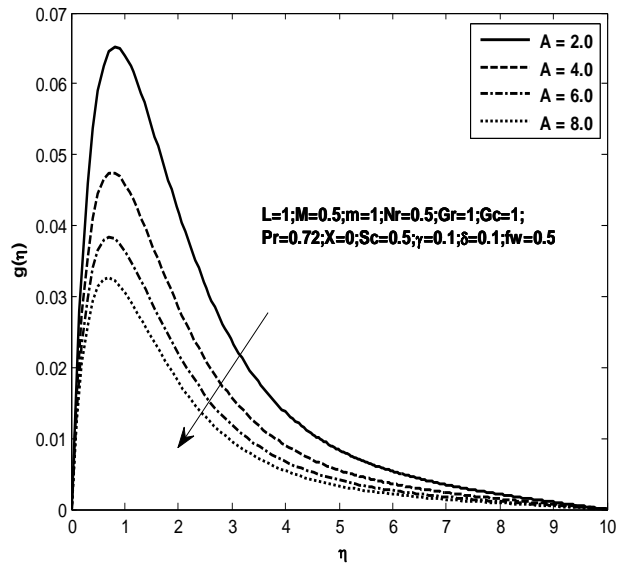


Fig. 6: Secondary velocity profiles for different values of A

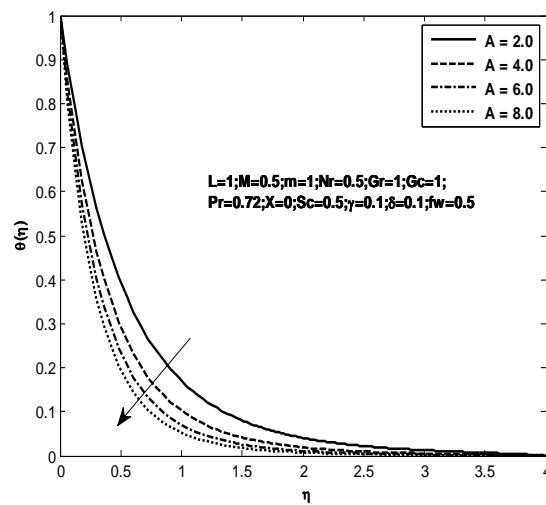


Fig. 7: Temperature profiles for different values of A

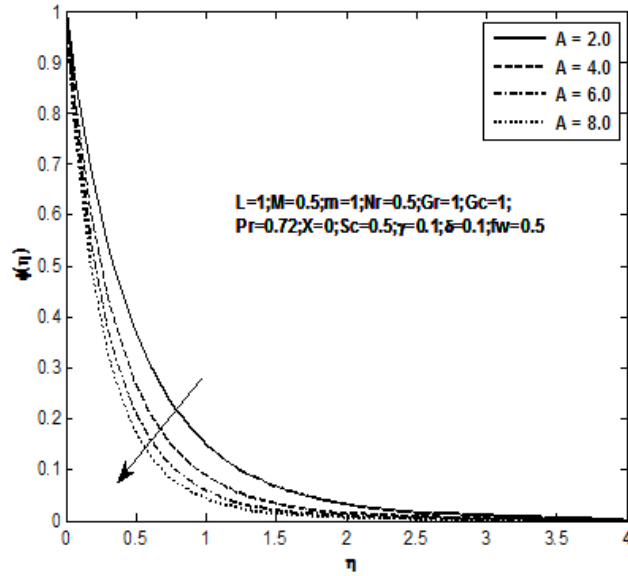


Fig. 8: Concentration profiles for different values of A

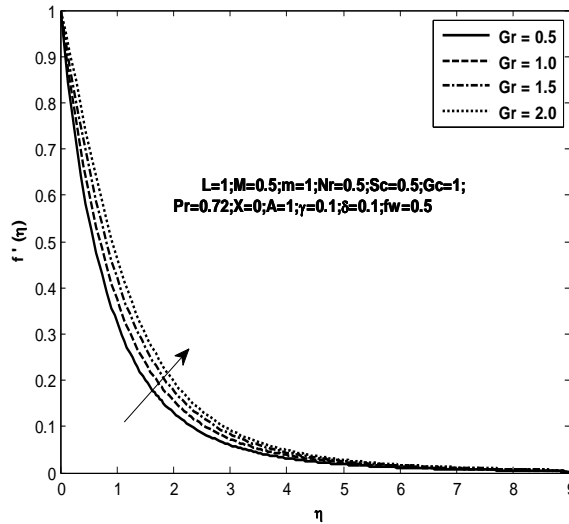


Fig. 9: Primary velocity profiles for different values of Gr

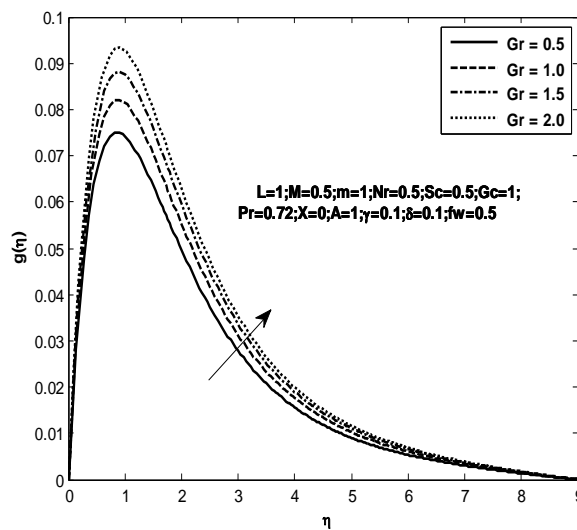


Fig. 10: Secondary velocity profiles for different values of Gr

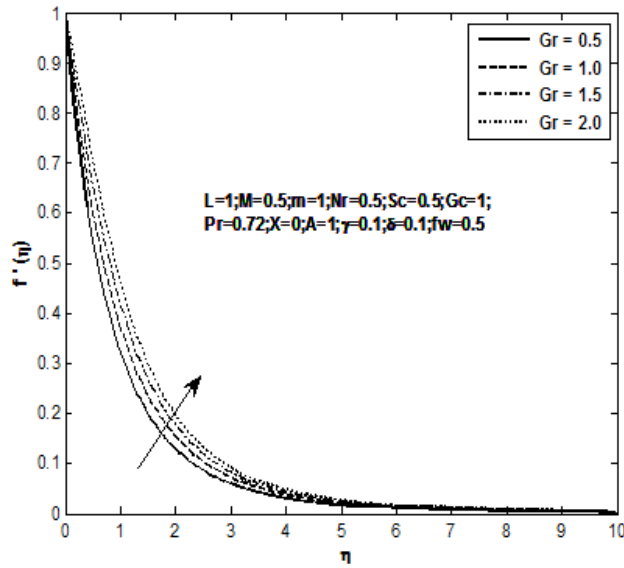


Fig. 11: Primary velocity profiles for different values of Gr

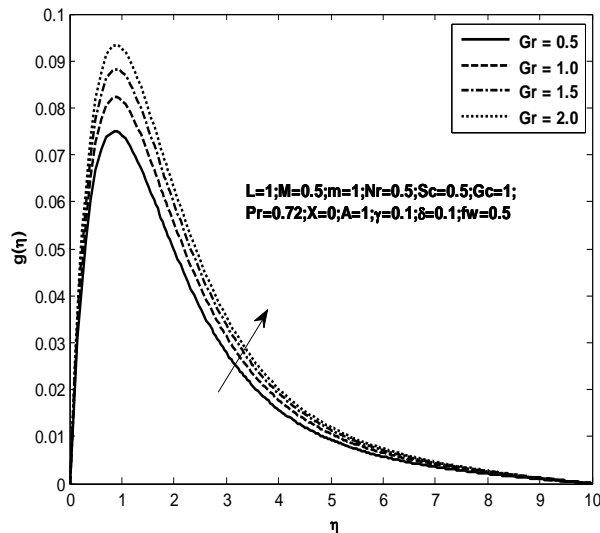


Fig. 12: Secondary velocity profiles for different values of Gr

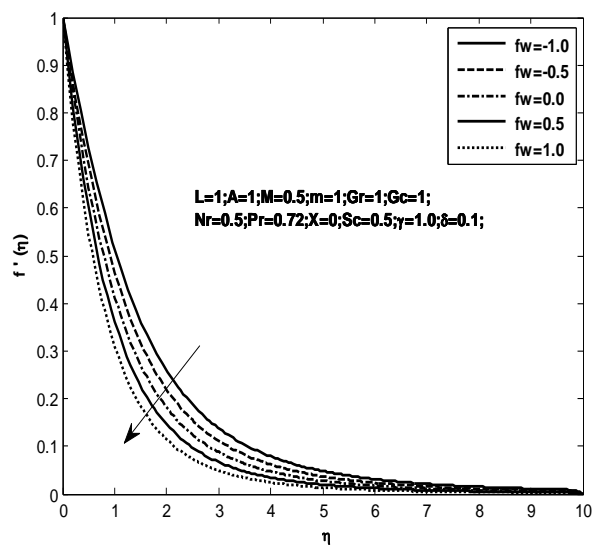


Fig. 13: Primary velocity profiles for different values of fw

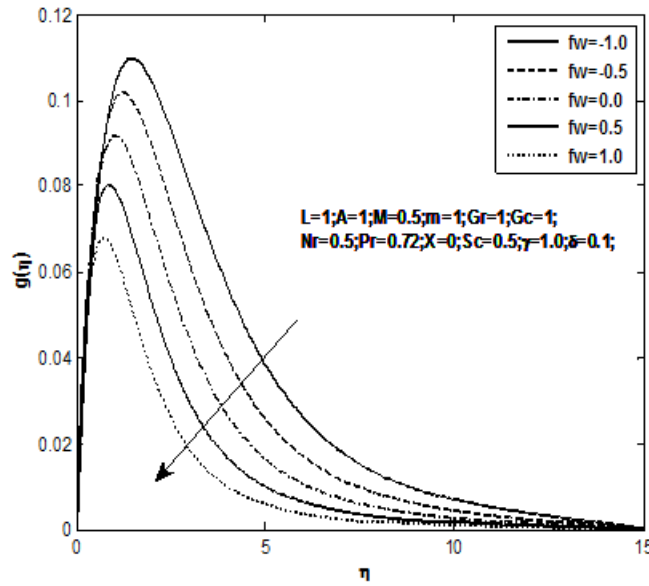


Fig. 14: Secondary velocity profiles for different values of f_w

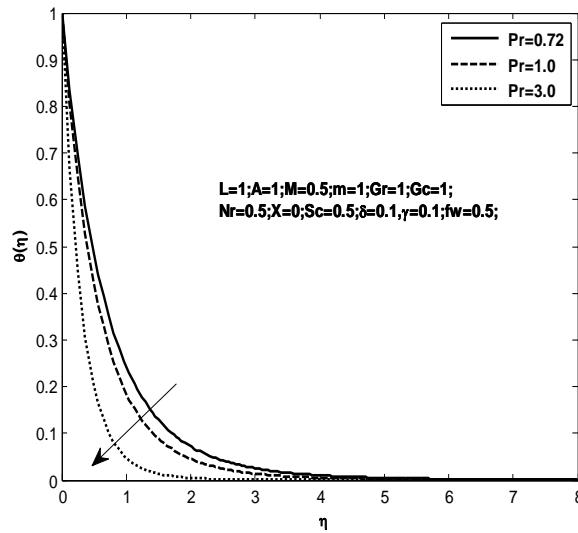


Fig. 15: Temperature profiles for different values of Pr

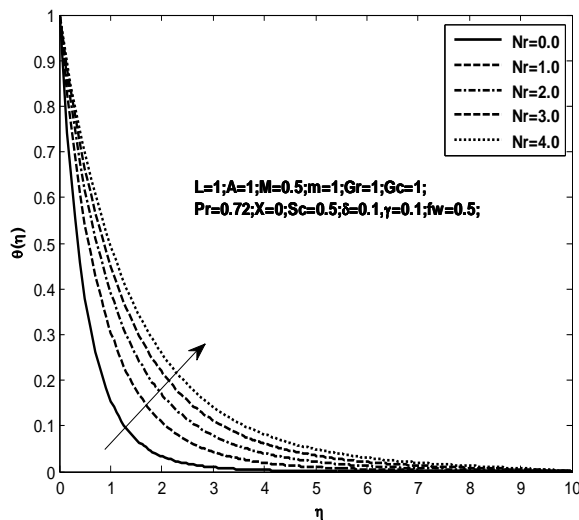


Fig. 16: Temperature profiles for different values of Nr

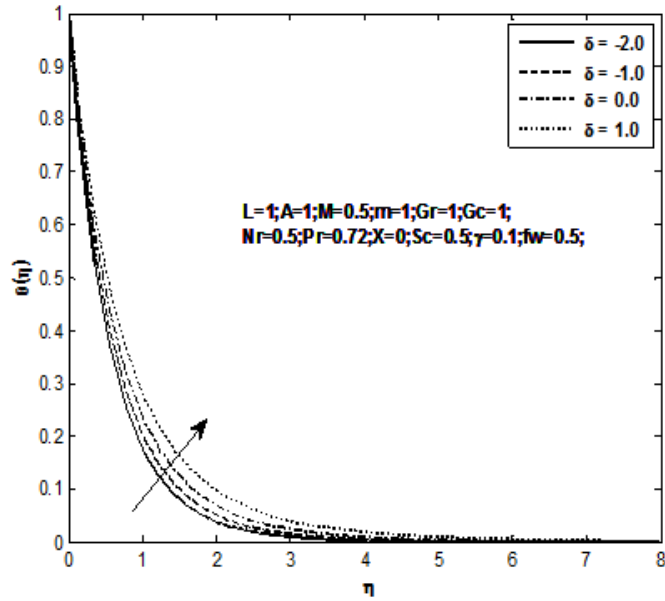


Fig. 17: Temperature profiles for different values of δ

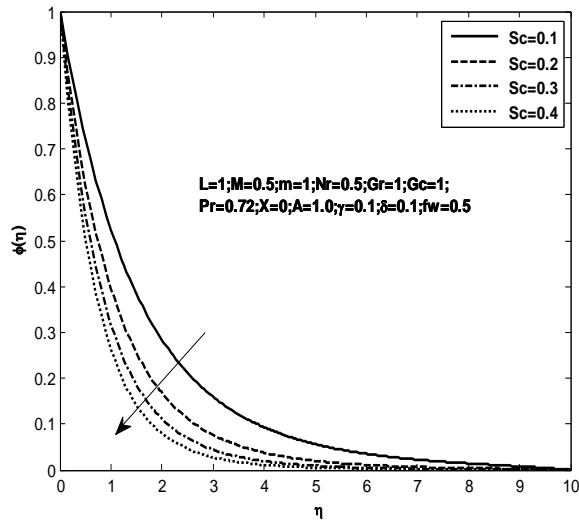


Fig. 18: Concentration profiles for different values of Sc

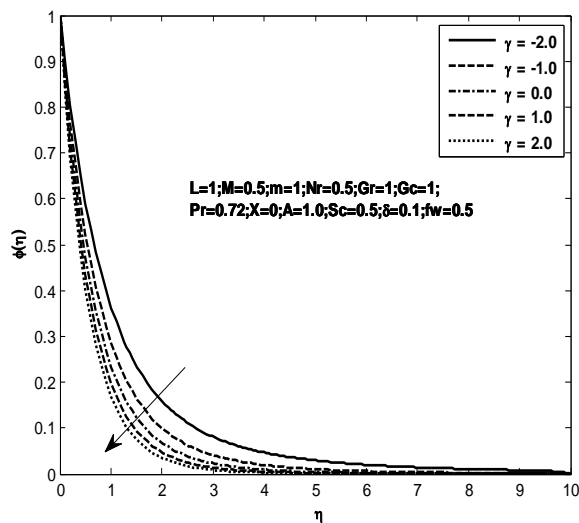


Fig. 19: Concentration profiles for different values of γ

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